

## The Crystal Structure of $\text{Rb}_5\text{Zr}_4\text{F}_{21}$

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The complex fluoride  $\text{Rb}_5\text{Zr}_4\text{F}_{21}$  crystallizes with space group  $P2_1$ ;  $a_0 = 11.520$  (5),  $b_0 = 11.222$  (5),  $c_0 = 7.868$  (2) Å and  $\cos \beta = -0.1445$  (3). Calculated density = 3.930 g.cm<sup>-3</sup> and  $Z = 2$ . The structure was solved by a tangent-formula procedure and was refined by Fourier and least-squares methods to a final discrepancy index ( $R$ ) of 0.0508 for 1376 observed  $\text{Cu K}\alpha_1$  reflections. The  $\text{F}^-$  coordination polyhedra are different for each of the four crystallographically independent  $\text{Zr}$  ions. Seven  $\text{F}^-$  ions are nearest neighbors to  $\text{Zr}(1)$  at the corners of a pentagonal bipyramid, and the interatomic distances are 1.99 (2) to 2.20 (2) Å. The second  $\text{Zr}$  ion has eight nearest neighbor  $\text{F}^-$  ions [1.98 (2) to 2.24 (2) Å] at the corners of an irregular antiprism.  $\text{Zr}(3)$  is octahedrally coordinated by 6  $\text{F}^-$  at distances of 1.90 (2) to 2.10 (2) Å. The fourth  $\text{Zr}$  ion has seven nearest neighbor  $\text{F}^-$  ions at the corners of an irregular antiprism with one corner missing. Interatomic distances  $\text{Zr}(4)-\text{F}$  are 1.95 (2) to 2.19 (2) Å. The  $\text{Rb}-\text{F}$  distances range from 2.68 (2) to 3.21 (3) Å. The structure is composed of cross-linked chains of  $\text{Zr}-\text{F}$  polyhedra. Chains are connected by  $\text{Zr}(2)-\text{Zr}(4)$  edge-sharing polyhedra and the space between the polyhedra is filled with  $\text{Rb}$  ions. Integrated intensities were measured by the  $\theta-2\theta$  scan technique using unfiltered  $\text{Cu K}\alpha$  radiation.

### Experimental

The compound  $\text{Rb}_5\text{Zr}_4\text{F}_{21}$  melts congruently at 445 °C (Thoma, 1959). Crystals picked from an ingot of the stoichiometric composition were ground to approximately spherical shape in an air-driven race. An ellipsoidal crystal of dimensions 0.26 × 0.364 × 0.2912 mm was mounted on a computer-controlled Picker four-circle goniostat equipped with a scintillation counter detector. Independent reflections  $hkl$ ,  $h$  and  $k \geq 0$  and all  $l$  were measured by the  $\theta-2\theta$  scan technique using unfiltered  $\text{Cu K}\alpha$  radiation. Each reflection was step-scanned in steps of 0.01°  $2\theta$  starting 0.75° before the calculated  $\text{Cu K}\alpha_1$  position and ending 0.75° beyond the calculated  $\text{Cu K}\alpha_2$  position. Each step was counted for 1 second and the background on each side of the peak was counted for 60 seconds. One reflection, 13 $\bar{1}$ , chosen as a standard reflection, was remeasured at intervals of 20 reflections. The net-count range of

13 $\bar{1}$  was 124700 to 121700. The 1403 independent reflections were corrected for Lorentz and polarization effects and absorption ( $\mu = 358.7$  cm<sup>-1</sup>). The maximum and minimum range of absorption correction was, respectively, 0.05626 and 0.00541. Conditions, for reflection  $0k0 = 2n$  and the diffraction symmetry  $2/m$ , are consistent with space groups  $P2_1$  and  $P2_1/m$ . A statistical analysis of the normalized structure factors established that the correct space group is  $P2_1$ , and this is confirmed by the structure. No piezoelectric tests were made on the crystals.

Unit-cell parameters were determined by a least-squares adjustment of 12 (80–100°  $2\theta$ )  $\text{Cu K}\beta$  reflections ( $\text{Cu K}\beta = 1.39217$  Å,  $T = 24^\circ\text{C}$ ). Cell parameters are;  $a_0 = 11.520$  (5),  $b_0 = 11.222$  (5), and  $c_0 = 7.868$  (2) Å;  $\cos \beta = -0.1445$  (3), ( $\beta = 98.3^\circ$ ). There are two formula weights per unit cell, and the calculated density is 3.930 gm.cm<sup>-3</sup>.

Initial positions of the heavy atoms were determined from a Fourier map plotted from normalized structure factors whose phases were determined by the tangent-formula method using program *PHASEM* (Drew,

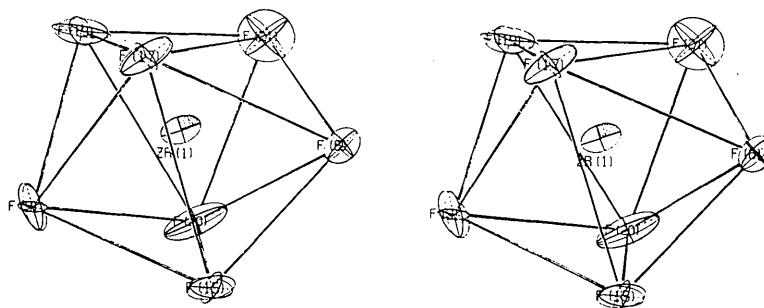


Fig. 1.  $\text{Zr}(1)$  coordination polyhedron.

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Table 1. Atomic parameters and temperature factors,  $\text{Rb}_5\text{Zr}_4\text{F}_{21}$ 

Standard errors in parentheses, corresponding to the last significant digit, are given by the variance-covariance matrix. Coefficients in the temperature factor of the form:

$$\exp [-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)].$$

Arbitrary value of  $y$  for Zr(1) to establish origin on  $z_1$  axis.

	$x$	$y$	$z$	$\beta_{11}$	$\beta_{22}$	$\beta_{33}$	$\beta_{12}$	$\beta_{13}$	$\beta_{23}$
ZR(1)	0.5932(2)	0.3000(0)	0.7830(3)	0.0043(2)	0.0026(2)	0.0044(4)	-0.0000(2)	0.0016(2)	0.0000(2)
ZR(2)	0.7721(2)	0.1589(3)	0.4964(3)	0.0036(2)	0.0028(2)	0.0046(4)	-0.0001(2)	0.0018(2)	-0.0001(2)
ZR(3)	0.0928(2)	0.1925(3)	0.8252(3)	0.0040(2)	0.0031(2)	0.0044(4)	0.0000(2)	0.0015(2)	0.0003(2)
ZR(4)	0.2758(2)	0.3401(3)	0.4668(3)	0.0039(2)	0.0029(2)	0.0046(4)	0.0002(2)	0.0018(2)	0.0004(2)
R8(1)	0.6091(3)	0.4758(4)	0.2818(4)	0.0051(2)	0.0070(4)	0.0076(5)	-0.0002(2)	0.0016(3)	0.0013(2)
R8(2)	0.4555(2)	0.1388(4)	0.1914(4)	0.0053(2)	0.0048(2)	0.0087(5)	0.0001(2)	0.0032(3)	-0.0002(3)
R8(3)	0.9057(3)	0.3282(4)	0.1650(4)	0.0059(2)	0.0067(3)	0.0093(6)	0.0006(2)	0.0034(3)	0.0012(3)
R8(4)	0.7455(3)	-0.0067(4)	0.0234(4)	0.0073(2)	0.0056(3)	0.0078(5)	0.0019(2)	0.0009(3)	-0.0005(2)
R8(5)	0.1152(3)	0.0194(4)	0.3139(4)	0.0061(2)	0.0044(3)	0.0098(6)	-0.0000(2)	0.0001(3)	0.0007(2)
F(1)	0.675(1)	-0.002(2)	0.390(2)	0.0051(1)	0.004(2)	0.0053(1)	-0.001(1)	0.0001(1)	0.001(2)
F(2)	0.612(2)	0.109(2)	-0.017(2)	0.009(2)	0.006(2)	0.008(2)	-0.002(1)	0.005(2)	0.002(2)
F(3)	0.234(2)	0.206(2)	-0.007(3)	0.009(2)	0.009(2)	0.018(5)	0.001(2)	-0.001(2)	-0.002(3)
F(4)	0.801(1)	0.005(2)	0.649(2)	0.007(1)	0.005(2)	0.006(3)	-0.002(1)	-0.000(2)	0.003(2)
F(5)	0.034(1)	0.348(2)	0.897(2)	0.007(1)	0.004(2)	0.011(3)	-0.000(1)	0.005(2)	-0.001(2)
F(6)	0.758(1)	0.221(2)	0.763(2)	0.007(1)	0.008(2)	0.007(3)	0.001(1)	0.002(1)	-0.002(2)
F(7)	0.154(2)	0.057(2)	0.715(2)	0.010(2)	0.004(2)	0.015(4)	0.002(1)	0.005(2)	-0.001(2)
F(8)	0.179(2)	0.306(2)	0.672(2)	0.009(2)	0.004(1)	0.010(3)	-0.000(1)	0.007(2)	-0.003(2)
F(9)	0.420(1)	0.294(2)	0.656(2)	0.002(1)	0.005(2)	0.011(3)	0.002(1)	-0.002(1)	0.002(2)
F(10)	0.690(1)	0.206(2)	0.260(2)	0.006(1)	0.005(2)	0.008(3)	0.006(1)	0.002(2)	0.004(2)
F(11)	0.953(1)	0.183(2)	0.635(2)	0.004(1)	0.007(2)	0.009(3)	-0.001(1)	0.002(2)	0.000(2)
F(12)	0.125(2)	0.293(2)	0.325(3)	0.005(1)	0.005(2)	0.024(5)	-0.002(1)	-0.006(2)	-0.001(3)
F(13)	0.382(1)	0.370(2)	0.298(2)	0.007(1)	0.003(1)	0.009(3)	-0.000(1)	0.005(2)	-0.002(2)
F(14)	0.883(1)	0.397(2)	0.347(2)	0.004(1)	0.007(2)	0.011(3)	-0.000(1)	0.003(2)	0.000(2)
F(15)	0.599(1)	0.199(2)	0.545(2)	0.005(1)	0.005(1)	0.006(3)	-0.000(1)	0.001(1)	-0.003(2)
F(16)	0.302(1)	0.168(2)	0.434(2)	0.007(1)	0.003(2)	0.010(3)	0.000(1)	0.002(2)	-0.001(2)
F(17)	0.560(1)	0.145(2)	0.888(2)	0.007(1)	0.003(1)	0.008(3)	0.001(1)	0.004(2)	0.002(2)
F(18)	0.810(1)	0.331(2)	0.481(2)	0.008(1)	0.003(2)	0.012(3)	-0.001(1)	0.003(2)	-0.001(2)
F(19)	0.497(2)	0.377(2)	0.944(2)	0.009(2)	0.010(2)	0.004(3)	0.000(2)	-0.001(2)	-0.001(2)
F(20)	0.604(2)	0.438(2)	0.625(2)	0.013(2)	0.002(1)	0.009(3)	0.000(1)	0.007(2)	0.001(2)
F(21)	0.712(2)	0.398(2)	0.944(3)	0.009(2)	0.005(2)	0.017(4)	-0.001(1)	0.001(2)	-0.004(2)

1967). Normalized structure factors were calculated by a least-squares method of Levy, Thiessen & Brown (1970). Positions of the fluorine atoms were determined from Fourier and difference Fourier maps. The structure was refined by iterative least-squares adjustment using a modification of computer program *ORFLS* (Busing, Martin & Levy, 1962). The quantity minimized by the least-squares program was  $\sum w|F_o^2 - |F_c^2|$  with weights,  $w$ , equal to the reciprocals of the variances estimated from the empirical equation:

$$\sigma^2(F_o^2) = \{T + B + [0.03(T - B)]^2\}/[A(Lp)^2],$$

where  $T$  = total counts,  $B$  = background counts,  $A$  = absorption correction, and  $Lp$  = Lorentz and polarization corrections (Brown & Levy, 1964). Scattering factors for the ions were taken from Cromer & Waber

(1965), and the anomalous dispersion terms for Rb and Zr for  $\text{Cu K}\alpha$  radiation were taken from Cromer (1965). Discrepancy indices were:

$$R^1 = \sum |F_o^2 - |F_c^2| / \sum |F_o^2| = 0.1014 \text{ for } 1376 \text{ reflections} > \sigma \\ = 0.1017 \text{ for all reflections.}$$

$$R^2 = \sum |F_o| - |F_c| / \sum |F_o| = 0.0508 \text{ for } 1376 \text{ reflections} > \sigma \\ = 0.0521 \text{ for all reflections.}$$

The standard deviation of an observation of unit weight  $[\sum w(F_o - F_c)^2 / (n_o - n_v)]^{1/2}$  is 2.637, where  $n_o$  is the number of reflections (1376) and  $n_v$  is the number of variables (271). Atomic parameters and temperature factors are listed in Table 1; interatomic distances are given in Table 2, and the observed and calculated structure factors are shown in Table 3. An extinction

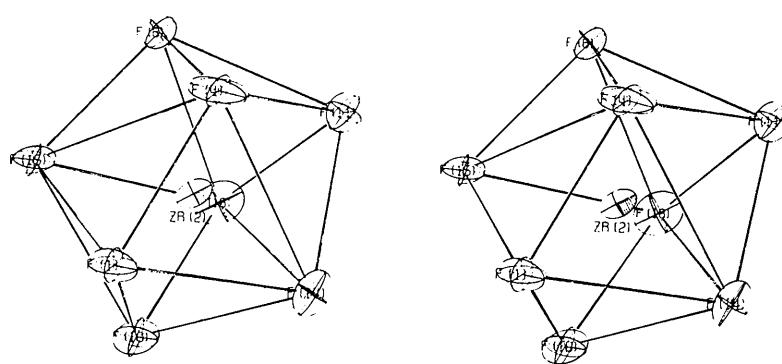


Fig. 2. Zr(2) coordination polyhedron.

Table 2. *Interatomic distances for Rb<sub>5</sub>Zr<sub>4</sub>F<sub>21</sub> in Å*

Zr(1)-F(1')	1.59(2)	Zr(1)-F(1'')	1.39(2)	Zr(1)-F(2)	2.00(2)
Zr(1)-F(21)	2.24(2)	Zr(1)-F(4)	2.10(2)	Zr(1)-F(1)	2.12(2)
Zr(1)-F(15)	2.20(2)	Zr(2)-F(14)	1.98(2)	Zr(2)-F(15)	1.99(2)
Zr(2)-F(10)	2.03(2)	Zr(2)-F(6)	2.10(2)	Zr(2)-F(1)	2.13(2)
Zr(2)-F(11)	2.22(2)	Zr(2)-F(11)	2.22(2)	Zr(2)-F(6)	2.24(2)
Zr(3)-F(2)	1.90(2)	Zr(3)-F(7)	1.93(2)	Zr(3)-F(3)	1.95(2)
Zr(3)-F(5)	1.93(2)	Zr(3)-F(11)	2.04(2)	Zr(3)-F(8)	2.10(2)
Zr(4)-F(13)	1.95(2)	Zr(4)-F(1)	1.95(2)	Zr(4)-F(12)	1.99(2)
Zr(4)-F(8)	2.14(2)	Zr(4)-F(5)	2.14(2)	Zr(4)-F(1)	2.14(2)
Zr(4)-F(4)	2.19(2)	Zr(5)-F(20)	2.74(2)	Zr(5)-F(7)	2.68(2)
Zr(5)-F(13)	2.90(2)	Zr(5)-F(17)	2.90(2)	Zr(5)-F(19)	3.00(2)
Rb(1)-F(16)	3.07(2)	Zr(5)-F(16)	3.17(2)	Rb(1)-F(10)	3.18(2)
Rb(1)-F(21)	3.19(2)	Zr(5)-F(10)	2.73(2)	Rb(2)-F(16)	2.80(2)
Rb(2)-F(20)	2.81(2)	Zr(5)-F(17)	2.52(2)	Rb(2)-F(3)	2.89(2)
Rb(2)-F(13)	2.90(2)	Zr(5)-F(15)	3.10(2)	Rb(2)-F(1)	3.19(2)
Rb(2)-F(19)	3.19(2)	Zr(5)-F(12)	2.56(2)	Rb(3)-F(1)	2.74(2)
Rb(3)-F(5)	2.55(2)	Zr(5)-F(7)	2.86(2)	Rb(3)-F(16)	2.86(2)
Rb(3)-F(14)	2.99(2)	Zr(5)-F(13)	3.02(2)	Rb(3)-F(2)	3.17(2)
Rb(4)-F(17)	2.84(2)	Zr(4)-F(5)	2.59(2)	Rb(4)-F(14)	3.04(2)
Rb(4)-F(13)	3.05(2)	Zr(4)-F(1)	3.11(2)	Rb(4)-F(4)	3.11(2)
Rb(4)-F(19)	3.12(2)	Zr(4)-F(13)	3.17(2)	Rb(4)-F(6)	3.19(2)
Rb(4)-F(3)	3.21(2)	Zr(5)-F(15)	2.72(2)	Rb(5)-F(16)	2.77(2)
Rb(5)-F(14)	2.65(2)	Zr(5)-F(2)	2.89(2)	Rb(5)-F(5)	2.93(2)
Rb(5)-F(12)	3.03(2)	Zr(5)-F(7)	3.15(2)	F(1)-F(10)	2.56(2)
F(1)-F(4)	2.32(2)	F(1)-F(9)	2.55(2)	F(1)-F(5)	2.82(2)
F(1)-F(14)	2.76(2)	F(1)-F(15)	2.76(3)	F(2)-F(5)	2.78(3)
F(1)-F(13)	3.00(3)	F(2)-F(3)	2.77(3)	F(3)-F(6)	2.75(3)
F(2)-F(11)	2.84(3)	F(2)-F(7)	2.91(3)	F(4)-F(12)	2.52(3)
F(3)-F(7)	2.80(3)	F(3)-F(5)	2.82(3)	F(4)-F(13)	2.57(3)
F(4)-F(5)	2.66(3)	F(4)-F(11)	2.67(3)	F(3)-F(2)	2.54(3)
F(4)-F(14)	2.57(3)	F(5)-F(2)	2.65(2)	F(3)-F(17)	2.75(2)
F(5)-F(11)	2.63(3)	F(6)-F(15)	2.34(2)	F(6)-F(12)	2.72(2)
F(6)-F(11)	2.52(2)	F(6)-F(12)	2.65(3)	F(6)-F(1)	2.95(3)
F(7)-F(11)	2.71(3)	F(7)-F(6)	2.83(3)	F(3)-F(15)	2.58(3)
F(8)-F(9)	2.50(3)	F(8)-F(11)	2.93(2)	F(1)-F(17)	2.60(3)
F(9)-F(19)	2.49(3)	F(9)-F(7)	2.50(3)	F(10)-F(14)	2.54(3)
F(9)-F(20)	2.71(3)	F(10)-F(1)	2.49(3)	F(11)-F(18)	2.52(2)
F(9)-F(13)	2.92(3)	F(11)-F(14)	2.49(3)	F(14)-F(16)	3.00(3)
F(10)-F(15)	2.61(2)	F(13)-F(15)	2.73(3)	F(15)-F(19)	2.95(3)
F(12)-F(16)	2.53(3)	F(15)-F(17)	2.57(3)	F(1)-(F(2))	2.45(3)
F(15)-F(20)	2.76(3)	F(17)-F(19)	2.45(3)		
F(16)-F(20)	2.85(3)				
F(20)-F(21)	2.67(3)				

correction was applied to  $F_c$  by Zachariasen's (1967, 1968) method. The range of extinction corrections was 0.7797–1.000 and Zachariasen's  $r^* = 0.0169$  (2) Å. The absolute configuration of the crystal could not be determined by Hamilton's (1965) significance test on the discrepancy indices. Ratio  $R_2(+\Delta f'')/R_2(-\Delta f'')$  for 1376 reflections  $>\sigma$  is 1.002, which is not significant at the 0.005 significance level for a one-dimensional hypothesis and 1105 degrees of freedom [ $R_2(-\Delta f'') = 0.0507$ ]. The same argument holds for any of the other discrepancy indices. Dr C. K. Johnson suggested that a similar test of absolute configuration could be made using a modification of ORFLS (Busing, Martin & Levy, 1962) which allows one to make  $\Delta f''$  a variable parameter in the least-squares refinement. All variable parameters of each of the two models,  $+\Delta f''$  and  $-\Delta f''$ , were refined first and subsequently  $\Delta f''$  was varied. In both cases, the least-squares refinement

diverges, and  $\Delta f''$  expands rapidly to a large positive number. The standard error for  $\Delta f''$  is 0.7 in both cases: cf.  $\Delta f''(\text{Rb}) = 1.81e$  and  $\Delta f''(\text{Zr}) = 2.42e$  for  $\text{Cu K}\alpha$  (Cromer, 1965).

Table 3. The observed and calculated structure factors, and phase angles for  $\text{Rb}_5\text{Zr}_4\text{F}_{21}$

**ALPHA** ≡ phase angle  $\alpha \times 10$

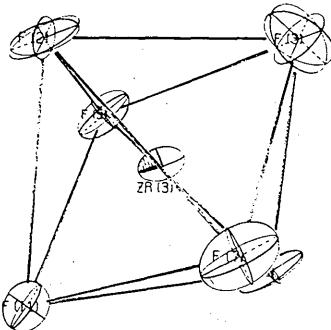
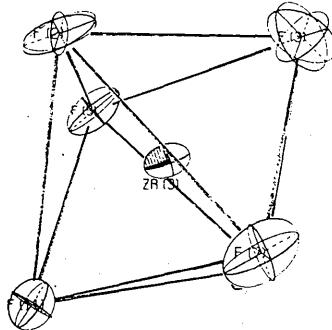


Fig. 3. Zr(3) coordination polyhedron.

## Results

The four Zr ions in the asymmetrical unit have seven, eight, six, and seven nearest neighbors, respectively, at the corners of irregular polyhedra illustrated in

Table 3 (cont.)

Figs. 1-4. The seven nearest neighbors to Zr(1) appear to be at the corners of a pentagonal bipyramid at distances 1.99(2) to 2.20(2) Å. Zr(2)'s nearest neighbors are at the corners of an irregular antiprism with interatomic distances of 1.98(2) to 2.24(2) Å. Octahedral coordination of Zr(3) is the easiest to see, and the distances are 1.90(2) to 2.10(2) Å. The seven nearest neighbors of Zr(4) are at the corners of a polyhedron that looks like an irregular antiprism with one corner missing. The Zr(4)-F distances are 1.95(2) to 2.19(2) Å. Interatomic distances are comparable to those of Burbank & Bensey (1956) (2.05 to 2.18 Å), Sears & Burns (1964) (2.05 to 2.16 Å) for an eight Zr-F coordination, and Brunton (1969) (2.012 to 2.167 Å) for a seven coordination. Interatomic distances for six-, seven- and eight-coordinated  $Zr^{4+}$  - F<sup>-</sup>, calculated from the empirical radii of Shannon & Prewitt (1969), are respectively 2.05, 2.11, and 2.13 Å for comparison.

The complete structure of  $\text{Rb}_5\text{Zr}_4\text{F}_{21}$  is shown in Fig. 5. Zr polyhedra, labeled according to the central occupant, form chains parallel to the  $a_0$  axis. In each chain of polyhedra, Zr(1) shares an edge with Zr(2) and a corner with Zr(4). Zr(2) also shares a corner with Zr(3). The chain is completed with Zr(3) sharing another corner with Zr(4). Chains are connected along the  $b$  axis by Zr(2) sharing an edge with Zr(4).

The Zr(1)–Zr(2) shared edge [F(6)–F(15), 2.34 (2) Å] and the Zr(2)–Zr(4) shared edge [F(1)–F(4), 2.32 (2) Å] are the shortest F–F distances in the structure (Table 2). The Zr(2)–Zr(3) corner is F(11) and the Zr(3)–Zr(4) corner is F(8). Zr–F distances are the greatest for the shared F atoms in each case for each Zr ion. The Rb ions fill the space between the Zr polyhedra.

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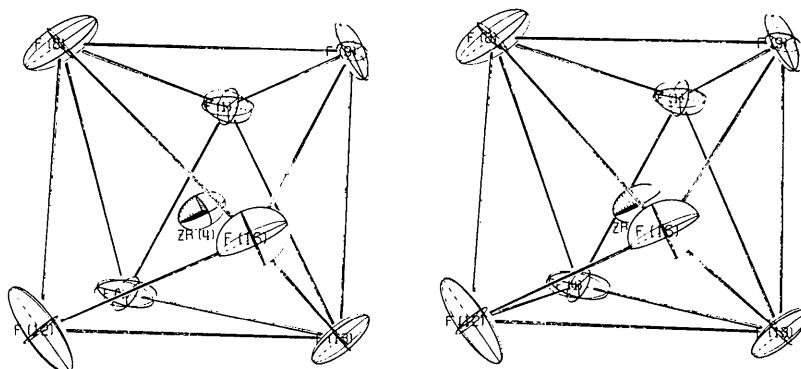
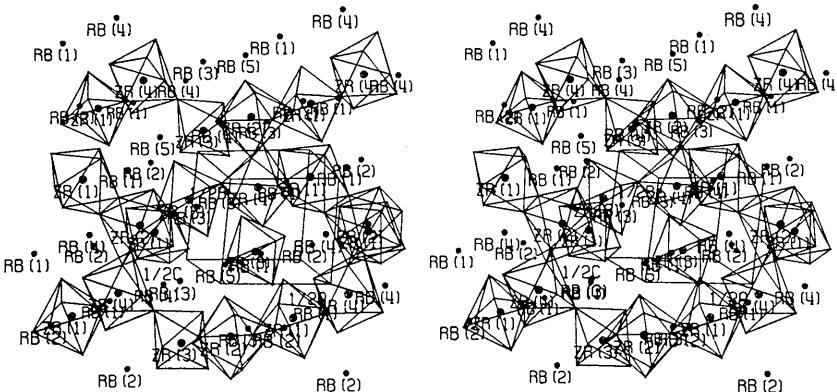


Fig. 4. Zr(4) coordination polyhedron.

Fig. 5. Structure of  $\text{Rb}_5\text{Zr}_4\text{F}_{21}$ .

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## The Crystal and Molecular Structures of Some Molecules Showing S···O Interaction. I. The Desaurin from Acetophenone.

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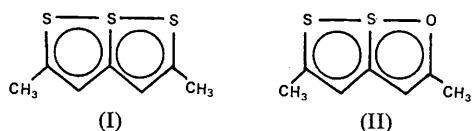
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The crystal structure of the 'desaurin from acetophenone',  $\text{C}_{18}\text{H}_{12}\text{O}_2\text{S}_2$ , has been determined as part of a program to investigate the S···O interactions in a range of molecules. The crystals belong to the monoclinic space group  $P2_1/c$  with  $a=5.740(8)$ ,  $b=4.848(12)$ ,  $c=28.299(8)$  Å,  $\beta=96.9(1)^\circ$ . The structure was solved from the Patterson function using 1020 photographically determined intensities and refined to  $R=0.124$ . The molecule shows an S···O distance of 2.640(6) Å.

### Introduction

The two-dimensional crystal structure analysis of 2,5-dimethylthiophthene (I) by Bezzi, Mammi & Garbuglio (1958)



indicated a symmetrical molecule with S···S distances of 2.36 Å. This early crystal structure analysis has since

been subjected to three-dimensional refinement in these Laboratories (Leung & Nyburg, 1969) and S···S determined as 2.358(1) Å. As pointed out by Bezzi *et al.* at the time, this is substantially longer than any 'formal' S-S bond known. A subsequent examination of S-S bond lengths by Hordvik (1966) suggests that the longest possible S-S bond, *i.e.* a single bond with no  $\pi$  character, would be 2.10 Å long. On the other hand, an examination of a number of sulphur-containing crystal structures shows that the minimum van der Waals radius of sulphur is about 1.65 Å [a value as high as 1.85 Å has been suggested by Pauling (1960)]. Accordingly, 2,5-dimethylthiophthene was the first